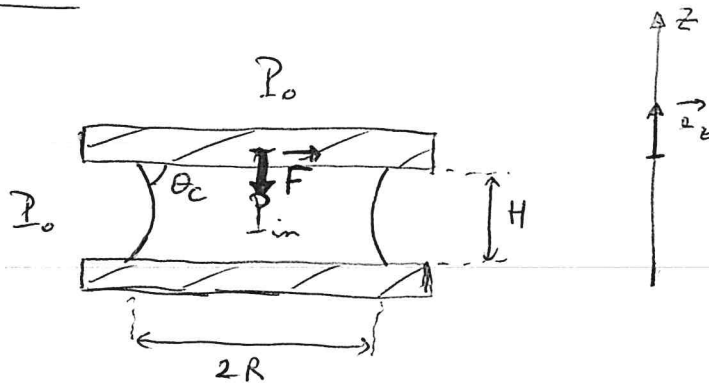


2/ Capillary adhesion



- $H \ll l_c \Rightarrow$ effects of gravity negligible.

- 2 contributions to the attractive force \vec{F} :

- capillary action = $-2\pi \delta R \sin \theta \vec{e}_z$

- pressure difference = $-\delta \mathcal{H} \pi R^2 \vec{e}_z$

where \mathcal{H} = mean curvature = $\delta \left(\frac{2 \cos \theta}{H} - \frac{1}{R} \right)$.

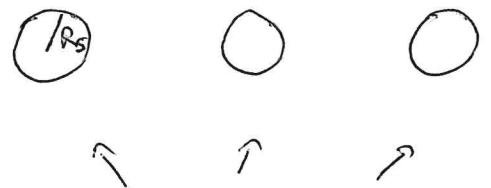
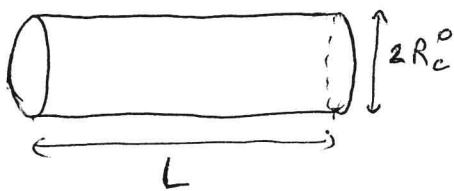
$$\Rightarrow \vec{F} = - \left[2\pi \delta R \sin \theta + \delta \left(\frac{2 \cos \theta}{H} - \frac{1}{R} \right) \pi R^2 \right] \vec{e}_z$$

when $H \ll R$:

$$\vec{F} = - \frac{2\delta \cos \theta}{H} \pi R^2 \vec{e}_z$$

4/ Rayleigh-Plateau instability

4.1 -



N spheres

Volume :

$$V = L \pi R_c^2$$

$$V = N \frac{4}{3} \pi R_s^3$$

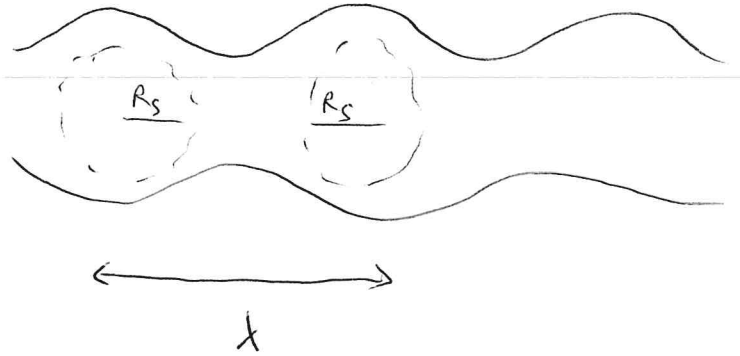
Surface :

$$S_{\text{cyl}} \approx L \times 2\pi R_c$$

$$S_{\text{sph}} = N \cdot 4\pi R_s^2$$

$$S_{\text{gh}} < S_{\text{cyl}} \quad \Leftrightarrow \quad \boxed{R_S > \frac{3}{2} R_c^0}$$

$$2/ \quad \lambda \simeq 2R_S \simeq 3R_c^0$$



$$4.2 - \quad \pi_c(x) = R_c + \delta\pi_c \cos(kx) \quad k = \frac{2\pi}{\lambda}$$

$$1. \quad \text{Volume conservat}^{\circ} : \int_0^{\lambda} \pi \pi_c^2(x) dx = \pi R_c^{\circ 2} \lambda$$

$$\Leftrightarrow \pi \left(R_c^2 + \frac{(\delta\pi_c)^2}{2} \right) \lambda = \pi R_c^{\circ 2} \lambda$$

$$\Leftrightarrow \boxed{R_c = \left[R_c^{\circ 2} - \frac{(\delta\pi_c)^2}{2} \right]^{1/2}}$$

$$2. \quad \delta\pi_c \ll R_c^{\circ} \Rightarrow \boxed{R_c \simeq R_c^{\circ} \left(1 - \frac{(\delta\pi_c)^2}{2R_c^{\circ 2}} \right)}$$

$$3. \quad \text{Energy variation (per period)} : \Delta E = \delta \int_0^{\lambda} 2\pi \pi_c(l) dl - \delta \int_0^{\lambda} 2\pi R_c^{\circ}(x) dx$$

(gravity neglected)

$$dl^2 = dx^2 + dr_c^2$$

$$\rightarrow dl \simeq dx \left(1 + \frac{1}{2} \left(\frac{dr_c}{dx} \right)^2 \right)$$

$$= dx \left(1 + \frac{(\delta\pi_c)^2}{2} k^2 \sin^2 kx \right)$$

$$\Rightarrow \Delta E = 2\pi\delta \int_0^\lambda \left[-\frac{(\delta n_c)^2}{4R_c^0} + \delta n_c \cos kx + R_c^0 \frac{(\delta n_c)^2}{2} k^2 \sin^2 kx \right] dx$$

$$\Delta E = 2\pi\delta \frac{\lambda (\delta n_c)^2}{4R_c^0} \left[R_c^0 k^2 \lambda - 1 \right]$$

$$4 - \Delta E < 0 \quad \Leftrightarrow \quad \lambda > 2\pi R_c^0 k^2$$

Note : the mode that will have the largest expanding factor is the one effectively selected among all the modes satisfying the above inequality.